TITLE:

ENHANCED MAGNETOCALORIC EFFECT MATERIAL

Inventors:

Laura J. Henderson Lewis 53 Oak Drive Calverton, New York 11933

Citizenship:

U.S.A.

This invention was made with Government support under contract number DE-AC02-98CH10886, awarded by the U.S. Department of Energy. The Government has certain rights in the invention.

BACKGROUND OF THE INVENTION

The present invention relates generally to materials that exhibit thermodynamically first-order magnetic transitions (i.e., magnetic state change accompaned by changes in the volume of the material) and, more particularly, to a method of enhancing the magnetocaloric effect (MCE) in such giant magnetocaloric materials.

10

The magnetocaloric effect (MCE) describes the adiabatic conversion of a magnetically-induced entropy change to the evolution or absorption of heat, with a corresponding rise or decrease in temperature. Figure 1 provides a schematic illustration of the magnetocaloric effect in a material 100. It is usual to quantify the MCE by the entropy change achieved by a change of magnetic field 102. The entropy change is determined from dc magnetic measurement by using the Maxwell relation:

$$\left(\frac{\partial S(T,H)}{\partial H}\right)_{T} = \left(\frac{\partial M(T,H)}{\partial T}\right)_{H}$$
[1]

which can then yield a magnetic entropy change ΔS_M :

$$\Delta S_M(T, \Delta H) = \int_{H_1}^{H_2} \left(\frac{\partial M(T, H)}{\partial T} \right)_H dH.$$
 [2]

In the previous expressions [1] and [2], T is the absolute temperature and H is the applied magnetic field. Typical (adiabatic) magnetocaloric temperature changes under an applied field change of 7 T range from $\Delta T_{ad} \approx 2.5$ K at T = 10 K for Nd to $\Delta T_{ad} \approx 12$ K at T \sim 180 K for Dy.

Systems employing the magnetocaloric effect are important for energy-efficient, low-CO₂ emission refrigeration, air conditioning for vehicles and buildings, as well as for sensitive temperature/heat sensor applications. Materials with a large magnetocaloric effect may also be utilized as heat pumps and, compared with the conventional vaporcycle refrigerator, the magnetic refrigerator is environmentally benign and has a number of advantages which include high efficiency, low mechanical vibration and compact size.

Briefly, an active magnetic regenerator (AMR) refrigerator employs a porous bed of a magnetic refrigerant working material that acts as both the coolant that produces refrigeration and the regenerator for the heat transfer fluid. As the magnetic working material is subjected to the application of a magnetic field, the particles of the material warm in an adiabatic manner from the MCE and absorb heat from the environment. As fluid flows through the particle bed from the cold end to the hot end, the working material warms the fluid via heat transfer. The heat from the fluid is removed at the hot heat sink in the heat exchanger. After the fluid flow is stopped, the magnetic field is removed

20

10

which then causes the magnetic working material to cool. The hot fluid is forced back to the now-cool porous bed of material where it is cooled by the bed. Remnant heat is removed from the fluid by the cold sink in the cold heat exchanger.

The potential applications of magnetic refrigeration are wide-ranging; with properly optimized performance, it is expected that they will be employed in building climate control, frozen food processing plants and supermarket chillers. Utilization of these materials can be envisioned in automotive and aircraft climate control, with an especially promising application of automotive climate control for zero-emission electric vehicles.

Magnetic refrigeration technology can accomplish those objectives in an environmentally-friendly manner, without the use of ozone-depleting chemicals such as CFCs (halogenated chlorofluorocarbons), HCFCs (hydrochlorofluorocarbons), HFCs (hydrofluorocarbons), PFCs (fluorocarbons) and SF₆ (sulfur hexafluoride), other hazardous chemicals (NH₃) and the production of additional greenhouse gases. The energy efficiency resulting from use of technologies employing magnetic refrigeration is anticipated to reduce the amount of energy consumed as well as reduce CO₂ emissions. Thus, two significant benefits of magnetic refrigeration technology are the replacement of CFC's (which will reduce the potential for global warming) and designing climate control in electric vehicles. The latter technology will allow a greater fraction of the available automobile power to be used for transportation rather than be exhausted for climate control.

20

10

With those advantages in mind, significant challenges to the technological development of these systems exist. Probably the most daunting technological hurdle remaining is the development of a cost-effective MCE material requiring practical magnetic fields and size considerations.

All magnetic materials, to a greater or lesser degree, exhibit a magnetocaloric effect. However, some materials, by virtue of a unique electronic structure or physical nanostructure, may display a significantly enhanced MCE, which may potentially be harnessed for technological application. In contrast to the MCE found in paramagnetic (PM) materials, the large MCE exhibited by ferromagnetic (FM) materials near their magnetic phase transition temperature (the Curie temperature T_C) renders them suitable as working materials for magnetic cooling at temperatures $T > \sim 20$ K, and up to the target temperatures appropriate for commercial, industrial and home refrigeration application and heat pump devices, 200 K - 400 K (approximately -70 °C ~ 130 °C).

10

20

Further enhancements of the ferromagnetic MCE are possible if the magnetic transition at T_C is accompanied by a crystallographic lattice distortion, as is often found in strongly-correlated electron systems. This lattice distortion may be either a contraction or an expansion of the atomic lattice, and may or may not include a change of symmetry. Such coupled magnetic-crystallographic changes are often referred to as first-order magnetic phase changes. The enhancement of the MCE arises because of the overall greater entropy change that occurs with a combined crystallographic and magnetic phase change. While materials such as amorphous and nanocrystalline alloys, intermetallic

compounds, perovskite-type oxides and, recently, pnictides and carbides, have all been topics of vigorous study over the past 5 years by virtue of their interesting MCE properties, by far the most promising MCE material is the intermetallic compound series based on the composition $Gd_5(Si_xGe_{1-x})_4$, where $0 \le x \le 0.5$.

The magnetocaloric effect produced by this family of compounds, heretofore referred to as GdSiGe, has been christened as "giant" because of its extremely large magnitude. In particular, the MCE of the GdSiGe alloys is reversible and exceeds that of any other known magnetic material by a factor of two. Another noteworthy characteristic of the GdSiGe family is that the Curie temperature, and hence the MCE, may be tuned with compositional variation. This feature allows the working temperature of the magnetic refrigerator to vary from 30 K to 276 K by adjusting the Si:Ge ratio. Moreover, the GdSiGe series of compositions is "metamagnetic", which means that the magnetic transition from ferromagnetic to paramagnetic behavior at the Curie point can be induced by applied field and pressure as well as increased temperature.

10

20

However, magnetocaloric materials made from gadolinium (Gd) and its alloys are generally very expensive and require a very large and, therefore, impractical magnetic fields on the order of 2 – 10T to yield a large magnetocaloric effect. For example, the typical AMR refrigerator described above, utilizing approx. 3 kg of Gd spheres, operates near room temperature in applied magnetic fields between 1.5T and 5T. This design provides a temperature span of 38 degrees for a field change of 5T, and it generates up to 600 W of cooling power in a 5 T field with an efficiency that approaches 60 % at 5 T.

While these may be impressive figures, a magnetic field of 5T can only be generated with a superconducting magnet that needs liquid helium to operate. Thus, the main drawback impeding the successful exploitation of the GdSiGe alloys is that the magnetic field magnitude required to obtain the spectacular magnetocaloric effect is simply too high for wide-spread commercial, home and transportation sector use.

U.S. Patent No. 5,743,095 to Gschneider, Jr. et al. discloses an improved Gd₅(Si_{1-x}Ge_x)₄ magnetic refrigerant that provides a high magnetocaloric effect and a high regenerator efficiency parameter. It is stated that the inclusion of a magnetically-soft alloying element, such as Mn, Fe, Co or Ni, to the Gd₅(Si_{1-x}Ge_x) compound optimizes the magnetocaloric effect properties of the refrigerant. However, the magnetic fields necessary to influence the magnetocaloric effect of this improved compound is still obtainable only with liquid-He-cooled superconducting magnets, which are not practical additions in typical applications.

10

20

Focusing solely on magnetoresistive materials, U.S. Patent No. 5,767,673 to Batlogg et al. describes an improved magnetoresistance obtained in a thin single crystal perovskite La_{2/3}Ca_{1/3}MnO₃ at extremely low fields when two magnetically-soft ferromagnetic (Mn,Zn)Fe₂O₄ bars were used in close proximity to the perovskite manganite. It is stated that the magnetically soft material can be thinly layered on the magnetoresistive core or it can be mixed with a magnetoresistive material to produce the improved magnetoresistive element. In each case, it is disclosed that the magnetically soft material placed in close proximity to the magnetoresistive core serves to increase the

magnetic field experienced by the magnetoresistive core resulting in an increased magnetoresistive effect at low applied magnetic fields.

While such advances have been made in the field of magnetoresistance, it is clear that innovative material design and engineering is needed to lower the applied magnetic field necessary to realize the optimum MCE in magnetocaloric materials. In particular, there is a great motivation to bridge the gap between giant magnetocaloric materials and the state-of-the-art AMR permanent magnet refrigerator design.

SUMMARY OF THE INVENTION

The present invention involves the application of a conformal thin layer of an elastically stiff material to a giant magnetocaloric material to effectively amplify the magnetic field applied to the magnetocaloric material and thereby lower the critical field necessary to realize a large MCE without adversely affecting the functionality of the material.

10

20

Thus, the present invention is a magnetocaloric effect heterostructure having a core layer of a giant magnetocaloric material and an elastically stiff material layer coated on at least one surface of the magnetocaloric material layer. The elastically stiff material layer restricts volume changes of the core layer during application of a magnetic field to the heterostructure.

Preferably, the core layer of magnetocaloric material is a compound of Gd₅(Si_{1-x} Ge_x)₄ and the elastically stiff coating layer is a low-coercivity, high-magnetization material selected from the group consisting of iron, cobalt, nickel and magnetic oxides.

Additionally, the elastically stiff material layer is preferably coated on opposite surfaces of the core layer and, more preferably, substantially encapsulates said core layer.

In a preferred embodiment, the present invention is in the form of a magnetocaloric effect composite powder including a plurality of core particles of a giant magnetocaloric material. Each of the core particles is encapsulated within a coating of elastically stiff material that restricts volume changes of the core particles during application of a magnetic field thereto.

Again, the material of the core particles is a compound of $Gd_5(Si_{1-x} Ge_x)_4$ and the elastically stiff coating is a low-coercivity, high-magnetization material selected from the group consisting of iron, cobalt, nickel and magnetic oxides. Also, the core particles are preferably substantially spherical and have a diameter of about 30 μ m, whereas the coating has a thickness between 50nm and 200nm.

10

20

The present invention further involves a method for enhancing the magnetocaloric effect within a giant magnetocaloric material including the processing step of coating a surface of the magnetocaloric material with an elastically stiff material. The elastically stiff material restricts volume changes of the magnetocaloric material during application of a magnetic field thereto.

In another method of enhancing the magnetocaloric effect within a giant magnetocaloric material according to the present invention, volume changes of a giant magnetocaloric material are restricted during application of a magnetic field thereto.

Preferably, the volume changes of the magnetocaloric material is restricted by a coating of an elastically stiff material.

As a result of the present invention, it is possible to significantly lower the applied magnetic fields necessary to obtain a large heating or cooling effect in giant magnetocaloric materials such as the $Gd_5(Si_{1-x}Ge_x)_4$ family of compounds. It is envisioned that this composite architecture can be easily deployed in many materials geometries and forms. The present invention can produce an effect that resembles amplification of an applied magnetic field and an increased or enhanced magnetocaloric effect.

Other objects, advantages and features of the present invention will become apparent from the following detailed description considered in conjunction with the accompanying drawings. It is to be understood, however, that the drawings are designed as an illustration only and not as a definition of the limits of the invention.

10

BRIEF DESCRIPTION OF THE DRAWINGS

Figure 1 is a schematic diagram illustrating the magnetocaloric effect (MCE) in materials.

Figure 2 is a cross-sectional view of an enhanced magnetocaloric effect material formed in accordance with the present invention.

Figure 3 is a cross-sectional view of the preferred embodiment of the enhanced magnetocaloric effect material of the present invention.

Figure 4 is a graph showing selected magnetization curves at three representative temperatures for prepared samples of the enhanced magnetocaloric effect material of the present invention.

Figure 5 is a graph showing the temperature dependence of the initial χ_{dc} of the samples shown in Figure 4.

Figure 6 is a graph showing the onset field differences ΔH_{onset} between the samples shown in Figure 4.

Figure 7 is a graph showing the entropy changes determined for internal field changes of 1 T and 5 T for the samples shown in Figure 4.

DETAILED DESCRIPTION OF THE PREFERRED EMBODIMENTS

10

20

In general, the present invention is a magnetocaloric effect (MCE) heterostructure having two different phase components. Figure 2 shows a cross-section of an enhanced MCE heterostructure material 10 of the present invention in its simplest form. The enhanced MCE heterostructure 10 includes a first component consisting of a giant magnetocaloric effect compound core 12. As mentioned above, a "giant magnetocaloric material" is a magnetic material that undergoes a coupled magnetic and structural change under the application of a magnetic field. The MCE core 12 of the present invention is preferably a compound of $Gd_5(Si_{1-x}Ge_x)_4$. The Gd-based core 12 is coated on at least one surface with a second component consisting of an elastically stiff material 14. An "elastically stiff" material, as defined herein, is a medium- to high-strength material that can restrict volume changes of the core 12 when a magnetic field is applied thereto. It has

been found that any material having a Modulus of Elasticity of more than about 0.1 x 10⁶ psi will meet this criteria. However, preferable elastically stiff materials will typically have a Modulus of Elasticity in the range of 10 x 10⁶ psi - 40 x 10⁶ psi and may include iron (Fe), cobalt (Co), copper (Cu), nickel (Ni), titanium (Ti), silicon (Si), aluminum (Al) and alloys thereof. Moreover, as will be discussed further below, it has been found that low-coercivity (magnetically soft), high-magnetization materials within this group, such as iron, cobalt, nickel and ferro- or ferrimagnetic magnetic oxides, are particularly well suited as a coating material 14 for the Gd-based compound core 12.

The magnetic field amplification coating 14 may be applied to slices of an arcmelted $Gd_5(Si_{1.5}Ge_{2.5})$ core 12. In the embodiment shown in Figure 2, the coating 14 is present as two conformal thin layers, having an approximate thickness of 100 nm (0.1 μ m), coated on opposite sides of the Gd-based core 12, whereas the Gd-based core has an approximate thickness of 0.6 mm. Thus, in this manner, a magnetic nanocomposite is created.

10

20

In the case where the coating 14 is a magnetically soft material, both components have distinct magnetocaloric properties at their respective ordering temperatures. As mentioned above, the Gd-based compound core alone exhibits a "giant" magnetocaloric effect around its magnetic ordering transition due to a first-order coupled magnetic-crystallographic transformation driven by the reversible destruction and reformation of specific covalent Si(Ge)-Si(Ge) bonds. The crystallographic transformation occurs via a martensitic-like collective shear movement of sub-nanometer thick building slabs. The

family of Gd₅(Si_{1-x}Ge_x) compounds is metamagnetic (i.e., the crystal structure transformation from the ferromagnetic orthorhombic Gd₅Si₄-type (*Pnma*) structure to paramagnetic monoclinic Gd₅Si₂Ge₂-type (*P*1121/*a*) structure can be easily induced by changes of temperature, magnetic field and pressure in the vicinity of the Curie temperature). While the metamagnetic character is reversible, it does have a temperature and field hysteresis consistent with the first-order thermodynamic nature of the phase change. Gd₅Si₁₋₅Ge_{2.5} is a magnetically heterogeneous system in the region of its first-order phase transition temperature characterized by coexistence of a magnetically-ordered (ferromagnetic) and a disordered (paramagnetic) phase. This two-phase coexistence gives rise to a unique two-stage initial magnetization behavior, wherein the first stage corresponds to magnetization of the ferromagnetic phase via domain-wall movement and magnetization rotation processes, and the second stage signals a field-induced paramagnetic-to-ferromagnetic (metamagnetic) transition that occurs at a critical onset field H_{onset}.

10

20

It has been found that the application of a thin layer of a strain- or deformation-resistant coating material 14 to both surfaces of a slice of Gd₅(Si₁₋₅Ge₂₋₅) core 12 effectively lowers the magnitude of the applied magnetic field required to induce the magnetocaloric effect within the core without significantly altering the magnitude of the induced MCE effect. The two phases of the resulting heterostructure complement each other to provide a "kick" to the MCE effect.

Figure 3 shows the preferred embodiment of a MCE nanocomposite 20 formed in accordance with the present invention. In order to more fully explore and optimize the results of the enhanced magnetocaloric effect investigations in magnetic nanocomposites it is preferred to fully encapsulate individual spherical particles 22 of the giant magnetocaloric material Gd₅(Si,Ge)₄ with a coating 24 of an elastically stiff material, as discussed above. The coating 24 is preferably conformal, homogeneous, uniform and of precise thickness. As also discussed above, soft magnetic materials are particularly well-suited as coating materials. The core particles 22 can be commercial-grade Gd₅(Si,Ge)₄ with nominal diameter of about 30 microns and a variety of coating compositions 24 and thicknesses can be applied. For example, coatings of 50 nm, 100 nm and 200 nm with compositions of Fe, Ni/Fe, Fe/Co, Al and Si can be applied to produce different samples. Thus, a MCE powder consisting of a plurality of nanocomposite particles 20 can be produced that can be utilized in a variety of applications.

10

20

The magnetic composite powder can be made by coating the GdSiGe particles 22 with an elastically stiff coating 24 using chemical vapor deposition. A Fast-Fluidized Bed Chemical Vapor Deposition (FFBCVD) process is ideal to achieve these processing parameters and objectives. Generally, chemical vapor deposition (CVD) is a plating method that relies on the chemical reaction of a vapor at a surface to form solid structural deposits. Since this is done on an atom-by-atom basis, impurity levels are typically less than 0.1% and densities are 99.9%. The CVD process utilizes a gaseous compound of the element to be deposited, which is then flowed over a heated substrate, resulting in thermal

decomposition or reduction of the gaseous compound and subsequent deposition of the material onto the substrate. The first layer forms at nucleation sites and, after the substrate is fully coated, growth continues on the deposit's crystal faces.

Successful CVD—producing dense, adherent coatings—depends upon experimentally determining optimal deposition parameters. These parameters include: the gaseous compound of the material to be deposited, substrate temperature, gas concentration, flow, pressure and geometry within the reaction chamber, coating thickness and substrate material. For the coating to have high integrity and adhesion to the substrate, the substrate either must have a similar coefficient of expansion to that of the deposited material or must form a strong chemical or metallurgical bond with it. The thinner the coating, the less similar the coefficients of expansion need be. Where coating and substrate form no intermetallic bond and have widely differing coefficients of expansion, a good bond often can be achieved by using a thin interlayer of a third material.

10

20

Fluidized beds, and more preferably fast-fluidized coating apparatus, are ideal tools for producing the fine powder materials of the present invention. Operating a fluidized bed above the transport limit (fast fluidization) or in the turbulent fluidization region may enable fine powders and whiskers, chopped fibers and such to be fluidized with high product yields. This technology eliminates problems with heat and mass transfer and handling fine and cohesive powders by using high gas shear velocities and turbulent gas mixing technologies. In a fast-fluidized bed, high gas velocities are used to

entrain the small particles, which are then filtered out of the gas entrainment stream and re-fed to the bottom of the reactor.

While not being bound by theory, it is believed that, in both the embodiments described above, both mechanical and magnetic mechanisms underlie the enhancement of the entropy change in the giant magnetocaloric compound Gd₅Si_{1.5}Ge_{2.5}. In particular, it has been found that the application of a ferromagnetic elastically stiff coating to the magnetocaloric compound has the effect of mechanically restricting the significant volume change that accompanies the magnetic transition occurring upon application of a magnetic field, thereby displacing the effect to different temperatures and fields. In magnetic terms, it is also believed that the stray field of a soft magnetic (e.g., Fe) layer serves to amplify the externally applied field. Measurements of the magnetostatic effect of the Fe coatings were carried out by comparing the room-temperature paramagnetic response obtained from the single-phase Gd₅Si_{1.5}Ge_{2.5} and the heterostructured Gd₅Si_{1.5}Ge_{2.5} specimens. No discernable difference in the response of the two specimens was detected, indicating that the magnetostatic contribution is extremely small and unlikely to produce the large changes noted in the ferromagnetic response. The enhanced field sensitivity found in the heterostructured samples is not due to a lowered energy barrier to reverse magnetic domain nucleation, because measured full field cycles indicate that the heterostructured samples display a stabilization of magnetization upon magnetic field reversal, not a reduction of the magnetization relative to the uncoated sample. Ferromagnetic exchange coupling between the Fe layer and the Gd₅Si_{1.5}Ge_{2.5} is

10

20

conceivable because of the very large value, approaching 50 nm, of the domain wall width in Fe despite the presence of a layer of Gd oxide that is undoubtedly intervening between the Fe and the Gd₅Si_{1.5}Ge_{2.5} interfaces.

It is envisioned that this composite architecture can be easily deployed in many materials geometries and forms. Established technology exists to provide a thin ferromagnetic metallic coating onto layers and particles of many compositions. Very important advantages can be envisioned with the implementation of this development. In the case of the giant magnetocaloric materials such as the $Gd_5(Si_{1-x}Ge_x)_4$ family of compounds these results strongly suggest that it is possible to significantly lower the applied magnetic fields necessary to obtain a large heating or cooling effect.

10

The following is an illustrative example demonstrating the foregoing:

Example

10

20

A $Gd_5Si_{1.5}Ge_{2.5}$ sample core slice was prepared by conventional arc melting procedures using 99.9 percent Lunex rare earths and 99.999 percent silicon and germanium. The weight loss after melting was less than 1 per cent and no subsequent heat treatment was carried out. Laboratory Cu- K_{α} x-ray diffraction studies revealed that the sample was well-crystallized and single-phase within the estimated limits of detection (5 vol%) and was isotropic, exhibiting no texture. Slices of $Gd_5Si_{1.5}Ge_{2.5}$ of approximate thickness 0.6 mm were cut with a water-cooled slow-speed wire saw and the surfaces were carefully polished with fine emery paper. Iron (Fe) layers of nominal thicknesses 0.1 μ m or 0.2 μ m were applied to both surfaces of the core slices by a chemical vapor deposition technique, with very small volume ratios of Fe: $Gd_5Si_{1.5}Ge_{2.5}$ of 0.03 vol% for the 0.1 μ m layers and was 0.06 vol% for the 0.2 μ m layers. During the Fe evaporation process the substrates were water-cooled and the chamber pressure was about 5×10^{-6} Torr. The deposition rate was ~3-4 Å/sec as monitored by a quartz crystal oscillator. Fig. 2 provides a schematic illustration of the heterostructured sample architecture.

Magnetization measurements were made on both an uncoated Gd₅Si_{1.5}Ge_{2.5} sample core and the coated heterostructure of Gd₅Si_{1.5}Ge_{2.5}+Fe in the temperature range 165 K to 235 K using a SQUID magnetometer. The uncoated Gd₅Si_{1.5}Ge_{2.5} specimen was cut into a single prism with dimensions 0.5mm x 0.5mm x 6.0 mm, while the coated slices were cut into uniform coupons of dimensions 1.26mm x 1.26mm x 0.6mm and

stacked to mimic the prism shape. Care was taken to determine the appropriate demagnetization factors of the specimens as accurately as possible. To this end, the specimens were experimentally corrected for demagnetization effects by using the demagnetization factor measured from Fe slices of the same dimensions and geometry as the specimens. An inherent assumption in the application of this procedure is that the stacked heterostructure slices act as a monolithic magnetic solid.

In the case of the heterostructure with the 0.1 μ m-thick Fe layer, ten slices were stacked together for measurement with the applied field perpendicular to the iron layer with a demagnetization factor $N_d = 0.083$. With the field applied parallel to the iron layers two slices were stacked together that produced a demagnetization factor of $N_d = 0.2475$. In the case of the heterostructure with 0.2 μ m Fe layer, two slices were stacked together for measurement with the applied field both perpendicular and parallel to the iron layer, and the demagnetization factors are calibrated to be 0.319 and 0.2475, respectively. The demagnetization factor of the single Gd₅Si_{1.5}Ge_{2.5} prism was determined to be $N_d = 0.022$. The estimated error of the calculated entropy change Δ S, based on consideration of the uncertainties in the measured parameters, is on the order of 2% in the vicinity of the zero-field Curie temperature.

10

20

Magnetization curves and ΔS determinations were made for a total of five sample configurations: the single-phase $Gd_5Si_{1.5}Ge_{2.5}$ prism and the heterostructured samples of 0.1- μ m and 0.2- μ m Fe thicknesses measured with the applied field oriented both parallel and perpendicular to the Fe film plane. Figure 4 shows selected magnetization curves at

three representative temperatures: at T = 188 K, below the zero-field Curie point; at the zero-field Curie point of 191 K and in the paramagnetic region at 215 K. For the sake of clarity only data for the parallel configuration is included in Figure 4 (the data obtained from samples in the perpendicular configuration are similar). The expected two-stage magnetization development of the samples in the magnetically-ordered state is evident.

It is noted that the Fe layers have basically no effect at T=215 K when the $Gd_5Si_{1.5}Ge_{2.5}$ alloy is in the paramagnetic state. However, in the magnetically-ordered state, in both parallel and the perpendicular sample geometries, the heterostructured sample data exhibit two main features that are attributed to the Fe layers. The first feature is that the magnetization development, as quantified by the magnitude of the initial dc magnetic susceptibility χ_{dc} , initiates at significantly lower fields in the heterostructured samples. Thus domain wall movement and rotation of magnetization in the ferromagnetic component of the $Gd_5Si_{1.5}Ge_{2.5}$ sample is promoted by the iron layer.

10

20

Figure 5 displays the temperature dependence of the initial χ_{dc} of the samples, where it can be seen that the Fe layers cause an increase of the initial susceptibility in the ferromagnetic component of the $Gd_5Si_{1.5}Ge_{2.5}$ alloy, effectively allowing it to be magnetized more readily. This susceptibility enhancement is especially evident at the lower temperatures and decreases monotonically with increasing temperature.

The second feature originating from the presence of the Fe layers is that the field-induced paramagnetic to ferromagnetic phase transition shifts to lower onset fields H_{onset} . The paramagnetic-ferromagnetic onset field is determined as the minimum of the field

derivative of the magnetization curve, dM/dH_i . The minimum in the curve signals the critical onset field at which the paramagnetic phase begins to change into the ferromagnetic phase, increasing the magnetic susceptibility in that field range. The onset field difference ΔH_{onset} between the single-phase and the heterostructured $Gd_5Si_{1.5}Ge_{2.5}$ data reaches a maximum at approximately 4500 G at 187 K, then decreases with increasing temperature as shown in Figure 6. The thicker Fe layer shifts the onset field to lower field values.

The increased field sensitivity of the magnetization processes and the metamagnetic transitions in the $Gd_5Si_{1.5}Ge_{2.5}$ compound conferred by the thin ferromagnetic coating results in an enhancement of the magnetocaloric effect. Figure 7 displays the entropy changes determined for internal field changes of 1 T and 5 T. It can be seen in Figure 7 that the selected heterostructured samples display a larger magnetic entropy change than the uncoated $Gd_5Si_{1.5}Ge_{2.5}$ in the vicinity of the metamagnetic transition temperature. For the sake of clarity of presentation, only the data from samples measured with the field parallel to the 0.1 μ m-thick Fe layers are shown (the entropy changes for the same samples measured with the field perpendicular to the Fe layers fall between the uncoated $Gd_5Si_{1.5}Ge_{2.5}$ data and the Fe-coated parallel data).

10

20

The maximum effect is the 11% enhancement achieved at $H_{int} = 5$ T with the field applied perpendicular to the 0.1 μ m Fe layers. Although thicker iron layers shift the metamagnetic transition to lower field values, they do not contribute to an enhanced magnetocaloric effect. This result indicates that magnetocaloric effect enhancement is

sensitive to the thickness of the Fe layer, and presents the possibility of obtaining an even more amplified MCE from optimized Fe layer thicknesses.

Thus, it has been found that in the metamagnetic transition temperature range the Fe layer situated in a parallel geometry lowers the onset field of the metamagnetic transition by 4240 Oe for a 0.1 µm layer thickness and by 4940 Oe for a 0.2 µm layer thickness from the base value of 8700 Oe at 187 K in the absence of Fe layers. Furthermore, the 0.1 µm layer of Fe oriented perpendicular to the applied field is found to enhance the entropy change, and thus the magnetocaloric effect, by approximately 11% above its base value at 191.5 K.

These results conclusively demonstrate that very thin layers of Fe can produce an effect that resembles amplification of the applied magnetic field. The effects of Fe layers with thicknesses of 0.1 μ m and 0.2 μ m are demonstrated by an increased initial dc susceptibility χ_{dc} and a shift of the onset field H_{onset} to lower fields to induce the paramagnetic-ferromagnetic metamagnetic transition in the $Gd_5Si_{1.5}Ge_{2.5}$ compound. Furthermore, the 0.1 μ m coating produces an increased entropy change, indicative of an increased magnetocaloric effect, in the vicinity of the $Gd_5Si_{1.5}Ge_{2.5}$ Curie transition. All of these effects are observed when the applied field is oriented in geometries both parallel and perpendicular to the iron layers.

10

20

Thus, while there have been described what are presently believed to be the preferred embodiments of the invention, those skilled in the art will realize that changes and modifications may be made thereto without departing from the spirit of the invention,

and is intended to claim all such changes and modifications as fall within the true scope of the invention.